# Some Remarks on Kuzmin's Theorem for F-Expansions 

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In this paper a general Kuzmin theorem for a class of multidimensional F-expansions is given. This result gives the uniform rate at which a certain sequence of approximates converges to the density of the invariant measure associated with an $F$-expansion. Some metric theorems are also given. This work extends and corrects some earlier results. The Jacobi algorithm is included as an example.

## 1. Introduction

In a recent paper [6], one of the authors gave a general Kuzmin theorem for a class of multidimensional $F$-expansions. The theorem applies to the Jacobi algorithm only when $n=2$. This is due to the fact that $T B_{v+1}=B_{v}$, $\nu \geqslant 1$, is not satisfied for $n \geqslant 3$. This difficulty can easily be shown by an example based on the restriction of digits [2,3]. In the proof of the theorem the assumption that $B_{v} \cap A_{i} \neq \varnothing$ implies $B_{v} \subseteq A_{i}, v \geqslant 1$, is implicitly used. Also, in the proof of Theorem 2 of [6] an inequality is obtained incorrectly. The notation of [6] is retained and we utilize much of the work done there.

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## 2. Kuzmin's Theorem

For $F \in \mathscr{F}$, we partition $(0,1)_{F}^{n}$ by the sets $T^{\nu} B_{v}(x), \nu \geqslant 1, x \in(0,1)_{F}^{n}$. This partition is assumed countable and is denoted by $\left\{A_{i}\right\}_{i \geqslant 1}$. Next we define

$$
E_{i, \nu}=\left\{\langle k\rangle=\left(k^{(1)}, k^{(2)}, \ldots, k^{(\nu)}\right) \mid T^{\nu} B(k) \supset A_{i}\right\}
$$

We remark that, if $B_{\nu}(\langle k\rangle)$ is a proper cylinder of order $\nu$, then $\langle k\rangle \in E_{i, \nu}$. Our first lemma generalizes Lemma 1 of [6].

Lemma 1. Suppose $F \in \mathscr{F}$. Let $\Psi_{0}$ be given and define $\Psi_{\nu}, \nu \geqslant 1$, by

$$
\begin{equation*}
\Psi_{\nu}(x)=\sum_{k \in E_{i, 1}} \Psi_{\nu-1}\left(f_{k}(x)\right)\left|J_{f_{k}}(x)\right|, \quad x \in A_{i} \tag{1}
\end{equation*}
$$

Then

$$
\begin{equation*}
\Psi_{\nu}(x)=\sum_{k \in E_{i, \nu}} \Psi_{0}\left(f_{v}(x)\right)\left|J_{v}(x)\right|, \quad x \in A_{i} \tag{2}
\end{equation*}
$$

where $f_{v}(x)=\prod_{i=1}^{v} \circ f_{k^{(i)}}(x)$ and $J_{v}=J_{f_{v}}$.
Proof. Our proof is by induction. Formula (2) is clearly true for $\nu=1$. Assume that (2) holds for $v$. Then, by definition,

$$
\Psi_{v+1}(x)=\sum_{k \in E_{i, 1}} \sum_{E_{j, \nu}} \Psi_{0}\left(f_{\nu}\left(f_{k}(x)\right)\right)\left|J_{v}\left(f_{k}(x)\right)\right|\left|J_{f_{k}}(x)\right|
$$

$x \in A_{i}$, where $j$ in the second summation satisfies $f_{k}(x) \in A_{j}$.
Since $f_{k}(x) \in A_{j}$ and $E_{j, \nu}$ is the set of all $\left(k^{(1)}, k^{(2)}, \ldots, k^{(\nu)}\right)$ such that $T^{\nu} B_{\nu} \supset A_{j}$, we have that $\left(k^{(1)}, k^{(2)}, \ldots, k^{(\nu)}, k\right)$ runs over all sequences ending in $k$ such that $f_{k^{(1)}} \circ f_{k^{(2)}} \circ \cdots \circ f_{k^{\prime \prime}} \circ f_{k}(x)$ is defined. Therefore, we have $x \in T^{\nu+1} B\left(k^{(1)}, \ldots, k^{(\nu)}, k\right)$ which implies $T^{\nu+1} B\left(k^{(1)}, \ldots, k^{(\nu)}, k\right) \supset A_{i}$. Conversely, if $T^{\nu+1} B\left(k^{(1)}, \ldots, k^{(\nu+1)}\right) \supset A_{i}$ we have

$$
T B\left(k^{(\nu+1)}\right) \supset T^{\nu+1} B\left(k^{(1)}, \ldots, k^{(\nu+1)}\right) \supset A_{i} .
$$

Therefore $k^{(\nu+1)} \in E_{i, 1}$. We have thus shown that ( $k^{(1)}, \ldots, k^{(\nu)}, k$ ) runs over all sequences of length $\nu+1$ in $E_{i, \nu+1}$.

Lemma 2. Suppose $F \in \mathscr{F}$ and define $\left\{\Psi_{\nu}\right\}_{v>0}$ as in Lemma 1. Then, if $\left\{A_{i}\right\}$ is countable,

$$
\int_{(0,1)^{n}} \Psi_{\nu}(x) d x=\int_{(0,1)^{n}} \Psi_{0}(x) d x, \quad v \geqslant 1
$$

Proof. The proof proceeds as in Corollary 6.1 of [5].

We next state our main theorem. The proof uses the technique of the proof of Theorem 2 in [6] but the difficulties mentioned in the introduction are corrected here. In addition, we make clear the uniformity of convergence to $\rho$.

Theorem 1. Let $F \in \mathscr{F}$ satisfy conditions (c), ( $q$ ), and ( $L$ ), and $\sigma(\nu)=\sup _{x} \operatorname{diam} B_{\nu}(x) \rightarrow 0$ as $\nu \rightarrow \infty$. In addition, assume the partition $\left\{A_{i}\right\}$ countable and that there is an $N \geqslant 1$ such that $B_{\nu} \cap T^{\mu} B_{\mu} \neq \varnothing$ implies $B_{\nu} \subseteq T^{\mu} B_{\mu}$ for all $\mu \geqslant 1$ and $\nu \geqslant N$. We suppose there exist constants $A$ and $D$ such that

$$
\left|\left(\partial f_{v}\right)_{k} / \partial x_{j}\right| \leqslant A \quad\left(x \in A_{i}\right)
$$

uniformly in $\nu, k$, and $j$; and

$$
\left\|J_{v}(x)|-| J_{v}(y)\right\| \leqslant \operatorname{Dm}\left(B_{v}\right)\|x-y\| \quad\left(x, y \in B_{v}\right)
$$

uniformly in $\nu$. Let $\left\{\Psi_{\nu}\right\}_{\nu>1}$ be the sequence of functions recursively defined in Lemma 1 with $\Psi_{0}$ satisfying

$$
0<m \leqslant \Psi_{0} \leqslant M \quad \text { and } \quad\left|\Psi_{0}(x)-\Psi_{0}(y)\right| \leqslant N\|x-y\|
$$

Then

$$
\begin{equation*}
\left|\Psi_{\nu}(x)-a \rho(x)\right|<b \sigma(\nu) \tag{3}
\end{equation*}
$$

where $\rho$ is the density of the invariant measure for $F, a=\int \Psi_{0} d m$ and $b$ are constants.

Proof. As in the proof of Theorem 2 of [6] we have

$$
\begin{equation*}
\left|\Psi_{\nu}(x)-\Psi_{\nu}(y)\right| \leqslant C_{1}\|x-y\| \tag{4}
\end{equation*}
$$

for $x, y \in A_{i}$, and

$$
g_{0} \Psi_{\nu}(x)<\Psi_{\nu+\mu}(x)<G_{0} \Psi_{\nu}(x)
$$

uniformly in $x, \mu$, and $\nu$.
We define

$$
\phi_{\nu}(x)=\Psi_{\mu+\nu}(x)-g_{0} \Psi_{v}(x)
$$

and

$$
\zeta_{\nu}(x)=G_{0} \Psi_{\nu}(x)-\Psi_{\mu+\nu}(x)
$$

By Lemma 1 and condition (C)

$$
\phi_{\nu}(x) \geqslant C^{-1} \sum^{(i)} \phi_{0}\left(f_{\nu}(x)\right) m\left(B_{v}\right)
$$

We let $\mathscr{P}_{v}=\bigcup B_{v}$ where the union is over all proper cylinders of order $\nu$. There exists $y_{v}{ }^{\prime} \in B_{v} \subseteq(0,1)_{F}^{n}$ such that

$$
\int_{B_{v}} \phi_{0} \leqslant \phi_{0}\left(y_{v}^{\prime}\right) m\left(B_{v}\right) .
$$

Hence

$$
-\int_{\mathcal{F}_{\nu}} \phi_{0} \geqslant-\sum_{\langle k\rangle} \phi_{0}\left(y_{\nu}^{\prime}\right) m\left(B_{v}\right) .
$$

Then for $v \geqslant N$

$$
\begin{aligned}
\phi_{v}(x)-C^{-1} \int_{\mathscr{F}_{v}} \phi_{0}(y) d y & \geqslant C^{-1} \sum_{\langle k\rangle}\left\{\phi_{0}\left(f_{v}(x)\right)-\phi_{0}\left(y_{v}^{\prime}\right)\right\} m\left(B_{v}\right) \\
& \geqslant-C^{-1} C_{1}\left(1+g_{0}\right) \sigma(\nu) \sum_{\langle k\rangle} m\left(B_{v}\right) \\
& \geqslant-C_{2} \sigma(\nu) .
\end{aligned}
$$

Here the sums are taken over all proper cylinders of order $\nu$ and the second inequality is by (4) applied to $\phi_{0}=\Psi_{\mu}-g_{0} \phi_{0}$. For the second inequality we use the fact that $f_{v}(x)$ and $y_{v}{ }^{\prime}$ are contained in the same element $A_{i}$ of the partition. This follows from the assumption that $B_{\nu} \cap T^{\mu} B_{u} \neq \varnothing$ implies $B_{\nu} \subset T^{\mu} B_{\mu}$ for $\mu \geqslant 1$ and $\nu \geqslant N$. This was tacitly used in [5] and [6]. That is,

$$
\Psi_{\mu+\nu}(x)-g_{0} \Psi_{\nu}(x) \geqslant C^{-1} \int_{\mathscr{P}_{\nu}}\left(\Psi_{\mu}-g_{0} \Psi_{0}\right) d m-C_{2} \sigma(\nu)
$$

In the same manner we obtain

$$
G_{0} \Psi_{\nu}(x)-\Psi_{\mu+\imath}(x) \geqslant C^{-1} \int_{\mathscr{P}_{v}}\left(G_{0} \Psi_{0}-\Psi_{\mu}\right) d m+C_{3} \sigma(\nu)
$$

The expressions for $l_{i}$ and $l_{i}{ }^{1}$ in [6] are functions of $g_{0}$ which makes $\nu_{0}$ a function of $g_{0}$. Therefore, when we iterate, it is possible that $\nu_{0}\left(g_{1}\right)>\nu_{0}\left(g_{0}\right)$ and so on. From this observation it is clear that we cannot let $r \rightarrow \infty$ without $\nu \rightarrow \infty$, and we see that inequality (9) of [6] was obtained incorrectly. We now give another derivation of a similar inequality.
We define

$$
\begin{aligned}
& \alpha(\nu)=1-\left(M_{1} C\right)^{-1} \int_{\mathscr{P}_{\nu}} \Psi_{0} d m, \\
& \beta(\nu)=\left(M_{1} C\right)^{-1} \int_{\mathscr{P}_{v}} \Psi_{\mu} d m-\frac{C_{3} \sigma(\nu)}{m_{1}},
\end{aligned}
$$

and

$$
\delta(\nu)=\left(M_{1} C\right)^{-1} \int_{\mathscr{P}_{\nu}} \Psi_{u} d m+\frac{C_{3} \sigma(\nu)}{m_{1}}
$$

From the equations

$$
g_{r+1}=g_{r}+\left(M_{1} C\right)^{-1} \int_{\mathscr{O}_{\nu}}\left(\Psi_{\mu}-g_{r} \Psi_{0}\right) d m-\frac{C_{2} \sigma(\nu)}{m_{1}}
$$

and

$$
G_{r+1}=G_{r}-\left(M_{1} C\right)^{-1} \int_{\mathscr{P}_{\nu}}\left(G_{0} \Psi_{0}-\Psi_{\mu}\right) d m+\frac{C_{3} \sigma(\nu)}{m_{1}}
$$

we obtain the recursion relations

$$
g_{r+1}=g_{r} \alpha(\nu)+\beta(\nu)
$$

and

$$
G_{r+1}=G_{r} \alpha(\nu)+\delta(\nu)
$$

We have the following inequalities

$$
\begin{aligned}
& \alpha(\nu) \leqslant 1-m_{1} q\left(M_{1} C\right)^{-1}<1 \\
& \beta(\nu) \geqslant \frac{m_{1} q}{M_{1} C}-\frac{C_{2} \sigma(\nu)}{m_{1}}>0 \\
& \delta(\nu) \geqslant m_{1} q\left(M_{1} C\right)^{-1}>0
\end{aligned}
$$

for $\nu \geqslant \nu_{0}$. We have assumed, without loss of generality, that $C>1$ and have used the inequality $m\left(\mathscr{P}_{\nu}\right) \geqslant q$. It now follows that

$$
g(\nu)=\lim _{r \rightarrow \infty} g_{r}=Q(\nu, \mu)+O(\sigma(\nu))
$$

and

$$
G(\nu)=\lim _{r \rightarrow \infty} G_{r}=Q(\nu, \mu)+O(\sigma(\nu))
$$

where

$$
Q(\nu, \mu)=\left(\int_{\mathscr{P}_{\nu}} \Psi_{u} d m\right) / \int_{\mathscr{P}_{\nu}} \Psi_{0} d m
$$

From

$$
g_{r} \Psi_{\nu}(x) \leqslant \Psi_{\mu+\nu}(x) \leqslant G_{r} \Psi_{\nu}(x)
$$

we have

$$
g(\nu) \Psi_{\nu}(x) \leqslant \Psi_{u+\nu}(x) \leqslant G(\nu) \Psi_{\nu}(x)
$$

and hence

$$
\left|Q(\nu, \mu) \Psi_{\nu}(x)-\Psi_{\mu+\nu}(x)\right|=O(\sigma(\nu))
$$

Now we integrate over $(0,1)^{n}$ and, by Lemma 2, obtain

$$
|Q(\nu, \mu)-1| \leqslant O(\sigma(\nu))
$$

Thus

$$
\left|\Psi_{\nu}(x)-\Psi_{\mu+\nu}(x)\right|=O(\sigma(\nu))
$$

and the proof is completed as in Theorem 2 of [6].
The assumption of Theorem 1 that there is an $N \geqslant 1$ such that $B_{v} \cap T^{\mu} B_{\mu} \neq \varnothing$ implies $B_{v} \subseteq T^{\mu} B_{\mu}$ for all $\mu \geqslant 1$ and $\nu \geqslant N$ is implied by the following condition: There is an $N \geqslant 1$ such that $T B_{v+1}=B_{v}$ for $\nu \geqslant N$. To see this, take $B_{\mu}=\bigcup B_{\mu+\nu}$. Then $T^{\mu} B_{\mu}=\bigcup T^{\mu} B_{\mu+\nu}=\bigcup B_{v}$. Since distinct cylinders are disjoint we are done. This observation will be used in Section 4 in dealing with the Jacobi algorithm.

## 3. Applications

As a first application of Theorem 1 we generalize Theorem 6 of [4] and Theorem 6.2 of [5]. The rate of convergence is the same as given in [4].

Theorem 2 (Gauss). Let F satisfy the assumptions of Theorem 1. If $\mu$ denotes the invariant measure, then

$$
\begin{equation*}
\left|m\left(T^{-\nu}(E)\right)-\mu(E)\right|<b m(E) \sigma(\nu) \tag{5}
\end{equation*}
$$

for all Borel sets $E$.
Proof. Define $\Psi_{0}(x)=1$. By Theorem 1, we obtain

$$
\begin{equation*}
\left|\Psi_{\nu}(x)-\rho(x)\right|<b \sigma(\nu) \tag{6}
\end{equation*}
$$

since $a=\int \Psi_{0}=1$. We multiply (6) by $I_{E}$ and integrate to obtain

$$
\left|\int I_{E} \Psi_{v} d m-\mu(E)\right|<b m(E) \sigma(\nu)
$$

Calculation similar to that in the proof of Theorem 6.2 in [4] shows

$$
\int I_{E} \Psi_{\nu} d m=m T^{-\nu} E
$$

and (5) follows.
The next application is a generalization of Theorem 7 of [4] and

Theorem 6.3 of [5]. We again obtain a rate of convergence identical with that in [4]. This result gives us a rate on the mixing.

Theorem 3. Let $F$ satisfy the assumptions of Theorem 1. Then for all Borel sets $F$ and cylinders $E$

$$
\begin{equation*}
\left|\mu\left(E \cap T^{-\nu} F\right)-\mu(E) \mu(F)\right| \leqslant \underset{q}{\langle C b\rangle} \mu(E) \mu(F) \sigma([\nu / 2]) \tag{7}
\end{equation*}
$$

Proof. First let $B_{v_{0}}$ be any proper cylinder and define

$$
\Psi_{0}(x)=\frac{I_{B_{\nu_{0}}}(x)}{\mu\left(B_{v_{0}}\right)} \rho(x)
$$

The sequence $\left\{\Psi_{\nu}\right\}_{v^{2} \nu_{0}}$ satisfies the hypotheses of Theorem 1. Our proof is completed as in Theorem 6.3 of [5], and we obtain

$$
\left|\mu\left(B_{\nu_{0}} \cap T^{-\nu-\nu_{0}}(F)\right)-\mu\left(B_{\nu_{0}}\right) \mu(F)\right| \leqslant \begin{gather*}
\langle C b\rangle  \tag{8}\\
q
\end{gather*} \mu\left(B_{\nu_{0}}\right) \mu(F) \sigma(v)
$$

A similar inequality holds for improper cylinders. The proof of inequality (7) from inequality (8) follows as in Folgerung 1 in [4].

## 4. The Jacobi Algorithm

The main purpose of this paper is to prove a Kuzmin theorem which applies to the Jacobi algorithm. In earlier statements of Theorem 1 [5], [6], it has been assumed that $\lim _{\nu \rightarrow \infty} \operatorname{diam} B_{v}(x)=0$ a.e. However, it is necessary to replace that assumption by $\sigma(\nu) \rightarrow 0$ as $\nu \rightarrow \infty$. In a paper of Paley and Ursell [1] it is shown, for $n=2$, that $\sigma(\nu)=O\left(\theta^{-\nu}\right)$, where $\theta$ is the unique root of $x^{3}-x^{2}-1=0$ with $1<\theta$. Unfortunately no proof is present for their results on the Jacobi algorithm with $n>2$. It is possible, however, to give a proof that $\sigma(\nu) \rightarrow 0$ as $\nu \rightarrow \infty$ which is valid for all $n \geqslant 1$. This proof, which will not be included here, is somewhat lengthly and depends on additional facts concerning the Jacobi algorithm.
An earlier paper [6] handled all other assumptions of Theorem 1 with the exception of proving $\left\{A_{i}\right\}$ countable and showing $T B_{v+1}=B_{v}$ for all $\nu \geqslant n-1$. Thus we must consider the partition generated by all sets $T^{\nu} B_{v}$.

We consider the transformation

$$
\begin{equation*}
T\left(x_{1}, x_{2}, \ldots, x_{n}\right)=\left(\frac{x_{2}}{x_{1}}-\left[\frac{x_{2}}{x_{1}}\right], \ldots, \frac{x_{n}}{x_{1}}-\left[\frac{x_{n}}{x_{1}}\right], \frac{1}{x_{1}}-\left[\frac{1}{x_{1}}\right]\right) \tag{9}
\end{equation*}
$$

which is defined on the unit cube after the deletion of an appropriate set of measure zero. This transformation corresponds to an $F$-expansion with

$$
F\left(x_{1}, \ldots, x_{n}\right)=\left(\frac{1}{x_{n}}, \frac{x_{1}}{x_{n}}, \ldots, \frac{x_{n-1}}{x_{n}}\right)
$$

We denote

$$
k^{(1)}(x)=\left(k_{1}^{(1)}, \ldots, k_{n}^{(1)}\right)=\left(\left[x_{2} / x_{1}\right], \ldots,\left[1 / x_{1}\right]\right)
$$

and

$$
k^{(s)}(x)=k^{(1)}\left(T^{s-1}(x)\right), \quad s \geqslant 1
$$

In this fashion we obtain a sequence of $n$-dimensional digits

$$
k^{(1)}, k^{(2)}, \ldots
$$

which is in one to one correspondence with the point $x$. The restrictions (see [1, 2]) on this sequence are as follows:

$$
\begin{equation*}
k_{n}^{(s)} \geqslant 1, \quad k_{n}^{(s)} \geqslant k_{i}^{(s)} \geqslant 0, \quad i=1,2, \ldots, n-1 \tag{10}
\end{equation*}
$$

and

$$
\begin{equation*}
\text { if } k_{i}^{(s)}=k_{n}^{(s)}, \quad \text { then } \quad k_{i-1}^{(s+1)} \leqslant k_{n-1}^{(s+1)} \tag{11}
\end{equation*}
$$

If we then obtain $k_{i-1}^{(s+1)}=k_{n-1}^{(s+1)}$, we have $k_{i-2}^{(s+2)} \leqslant k_{n-2}^{(s+2)}$, and so on. We formally set $k_{0}^{(t)}=1$ and $k_{r}^{(t)}=0$ for $r<0$.

The dependence of the digits is of length $n-1$. That is, the occurrence of a digit satisfying (10) depends only on the $n-1$ preceeding digits. Thus, we have $T B_{v+1}=B_{\nu}$ for $\nu \geqslant n-1$. The structure of $T^{-1}$, namely

$$
x_{1}=\frac{1}{k_{n}+y_{n}}, \quad x_{2}=\frac{k_{1}+y_{1}}{k_{n}+y_{n}}, \ldots, \quad x_{n}=\frac{k_{n-1}+y_{n-1}}{k_{n}+y_{n}}
$$

shows that the restrictions (11) correspond to ones of the type $y_{j}<y_{j+h}$. Therefore, the sets $T^{\nu} B_{\nu}$ partition the unit cube into $n!$ simplices $A_{i}$ defined by the intersections of the sets $\left\{\left(x_{1}, \ldots, x_{n}\right) \mid x_{j}<x_{j+h}\right\}$ and their complements. This completes our proof that Theorem 1 (and, hence, Theorems 2 and 3) applies to the Jacobi algorithm.

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